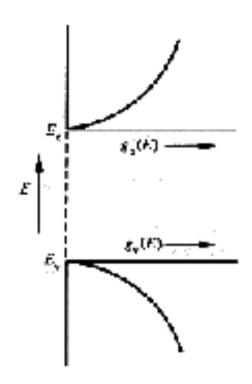


# Density of States and Fermi Energy Concepts



#### ■ Density of States Concept



$$g_c(E)dE$$

The number of conduction band states/cm<sup>3</sup> lying in the energy range between E and E + dE (if  $E \ge E_c$ ).

$$g_{v}(E)dE$$

The number of valence band states/cm<sup>3</sup> lying in the energy range between E and E + dE (if  $E \le E_v$ ).

General energy dependence of  $g_c(E)$  and  $g_v(E)$  near the band edges.



#### **Density of States Concept**

Quantum Mechanics tells us that the number of available states in a cm<sup>3</sup> per unit of energy, the density of states, is given by:

$$g_c(E) = \frac{m_n^* \sqrt{2m_n^*(E - E_c)}}{\pi^2 \hbar^3}, E \ge E_c$$
 Density of States in Conduction Band

$$g_{\nu}(E) = \frac{m_{\nu}^* \sqrt{2m_{\nu}^*(E_{\nu} - E)}}{\pi^2 \hbar^3}, E \leq E_{\nu}$$
 Density of States in Valence Band

in Valence Band

$$unit \equiv \left(\frac{Number of States}{cm^3}\right)_{eV}$$



- □ Probability of Occupation (*Fermi Function*) Concept
  - ➤ Now that we know the number of available states at each energy, then how do the electrons occupy these states?
  - ➤ We need to know how the electrons are "distributed in energy".
  - Again, Quantum Mechanics tells us that the electrons follow the "Fermi-distribution function".

$$f(E) = \frac{1}{1 + e^{(E - E_f)/kT}}$$

$$E_f \equiv \text{Fermi energy (average energy in the crystal)}$$

$$k \equiv \text{Boltzmann constant } (k = 8.617 \times 10^{-5} \text{eV/K})$$

$$T \equiv \text{Temperature in Kelvin (K)}$$

- $\Leftrightarrow$  f(E) is the probability that a state at energy E is <u>occupied</u>.
- 1-f(E) is the probability that a state at energy E is <u>unoccupied</u>.
- Fermi function applies only under <u>equilibrium conditions</u>, however, is universal in the sense that it <u>applies with all materials</u>-insulators, semiconductors, and metals.



#### ☐ Fermi-Dirac Distribution

- Applies to the particles that obey the exclusion principle. (Fermions)
- Fermi-Dirac Function: The probability that a particular state E is occupied for a systems of Fermions.

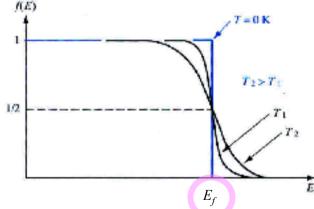
$$f(E) = \frac{1}{1 + \exp[(E - E_F)/kT]}$$

- E<sub>F</sub>: Fermi Energy (Fermi Level)
- Some properties of F-D distribution function:

$$f(E) \le 1$$

$$f(E = E_F) \le \frac{1}{2}$$

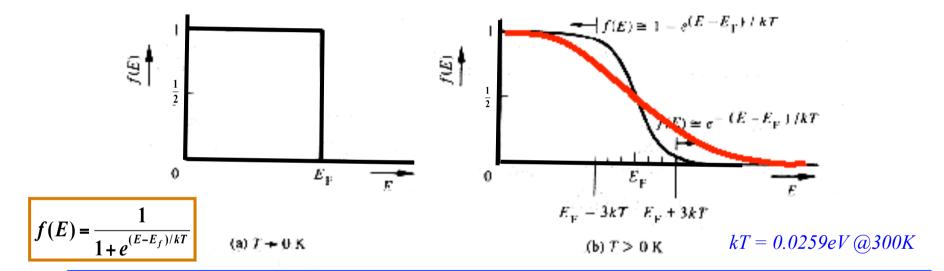
$$f(E_F + \delta E) = 1 - f(E_F - \delta E) = \frac{1}{1 + \exp(\delta E/kT)}$$



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### How do electrons and holes populate the bands?

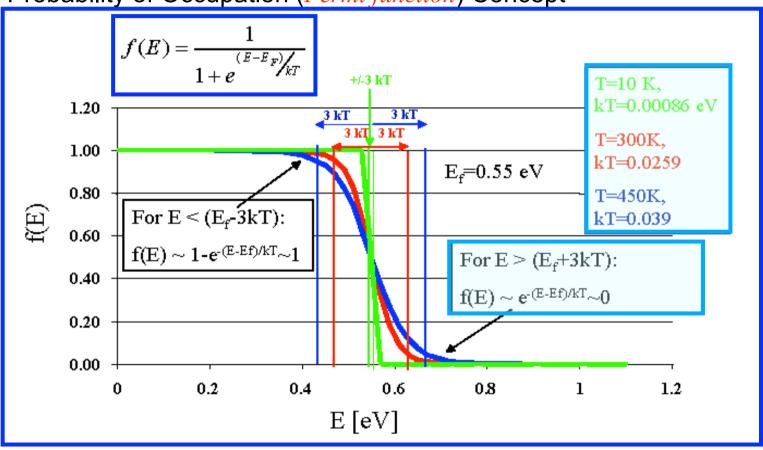
☐ Probability of Occupation (*Fermi function*) Concept



- At T=0K, occupancy is "digital": No occupation of states above  $E_f$  and complete occupation of states below  $E_f$ .
- At T>0K, occupation probability is reduced with increasing energy.  $f(E=E_f) = 1/2$  regardless of temperature.
- At higher temperatures, higher energy states can be occupied, leaving more lower energy states unoccupied  $[I f(E_f)]$ .



☐ Probability of Occupation (*Fermi function*) Concept



- $\blacktriangleright$  If  $E \ge E_f + 3kT \rightarrow e^{(E-E_f)/kT} \rangle \rangle$  1 and  $f(E) \approx e^{-(E-E_f)/kT}$
- ightharpoonup Consequently, above  $E_f + 3kT$  the Fermi function or filled-state probability decays exponentially to zero with increasing energy.



#### ☐ Probability of Occupation Concept

The density of electrons (or holes) occupying the states in energy between E and E + dE is:

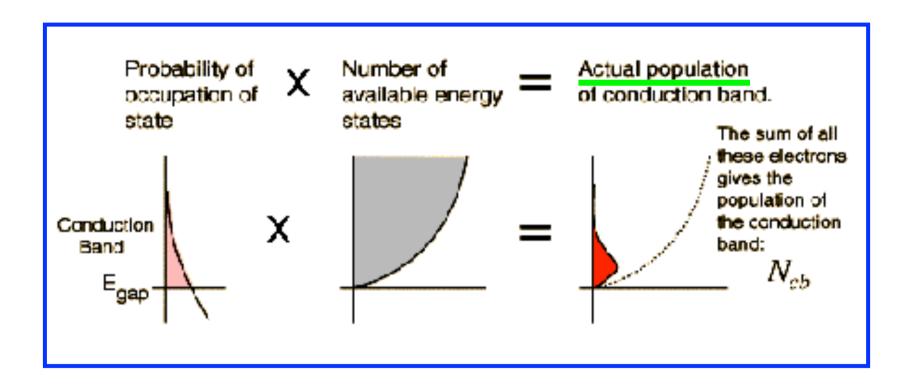
$$g_c(E)f(E)dE$$
 Electrons/cm<sup>3</sup> in the conduction band between  $E$  and  $E+dE$  (if  $E \ge E_c$ ).

$$g_v(E)f(E)dE \qquad \text{Holes/cm}^3 \text{ in the conduction band} \\ between $E$ and $E+dE$ \\ (if $E \le E_v$).$$

Otherwise

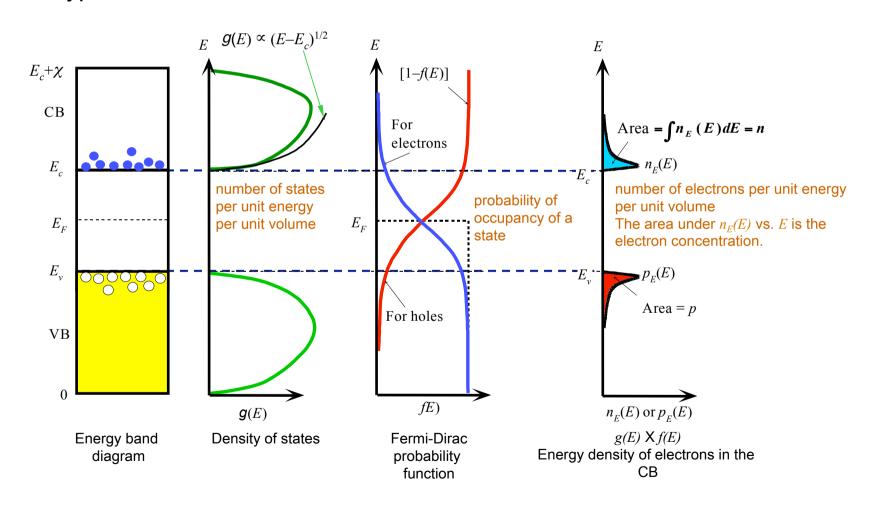


☐ Probability of Occupation Concept





#### ☐ Typical band structures of Semiconductor



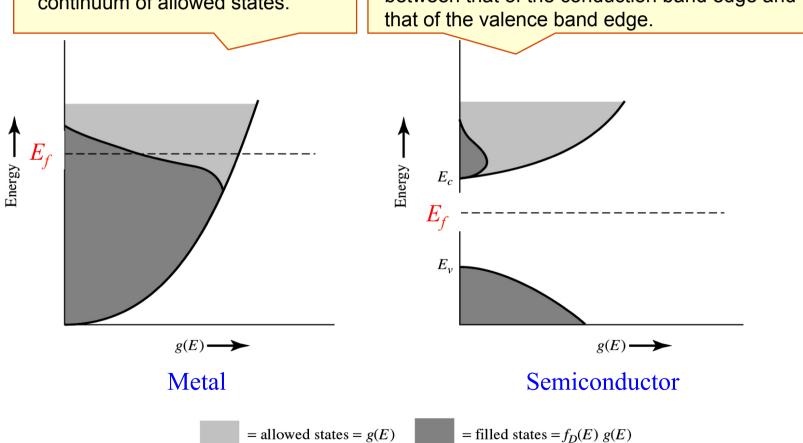


#### Metals vs. Semiconductors

 $\square$  Allowed electronic-energy states g(E)

Fermi level  $E_f$  immersed in the continuum of allowed states.

The Fermi level  $E_f$  is at an intermediate energy between that of the conduction band edge and that of the valence band edge.

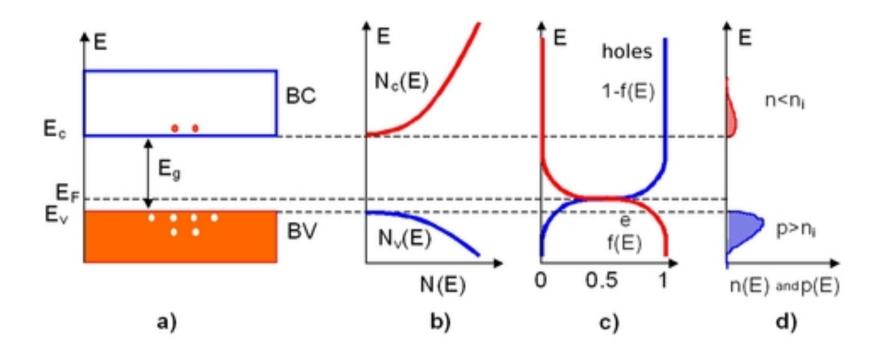




- ☐ Fermi function and Carrier Concentration
  - ➤ Note that although the Fermi function has a finite value in the gap, there is no electron population at those energies. (that's what you mean by a gap)
  - ➤ The <u>population</u> depends upon the <u>product of</u> the <u>Fermi function</u> and the <u>electron density of states</u>. So in the gap there are no electrons because the <u>density of states is zero</u>.
  - ➤ In the conduction band at 0K, there are no electrons even though there are plenty of available states, but the Fermi function is zero.
  - At high temperatures, both the density of states and the Fermi function have finite values in the conduction band, so there is a finite conducting population.

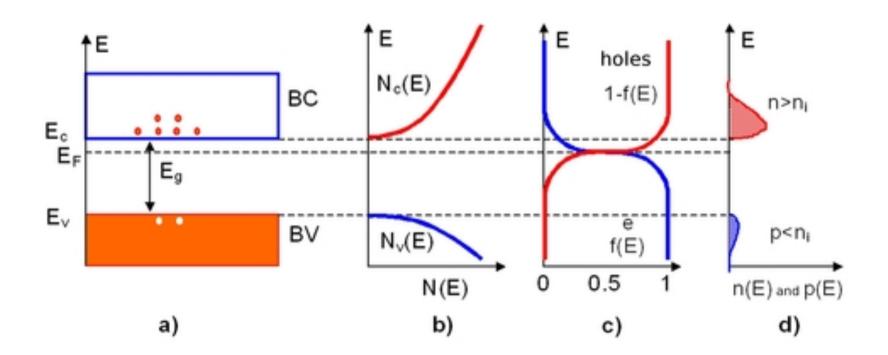


☐ Energy Band Occupation in p-type semiconductor



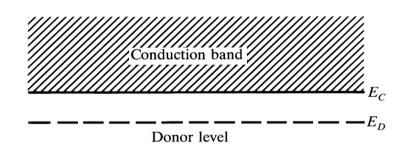


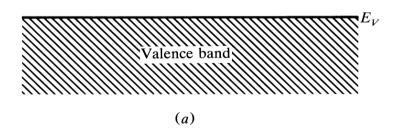
☐ Energy Band Occupation in n-type semiconductor

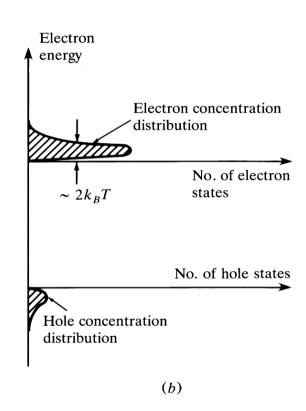




n-type material

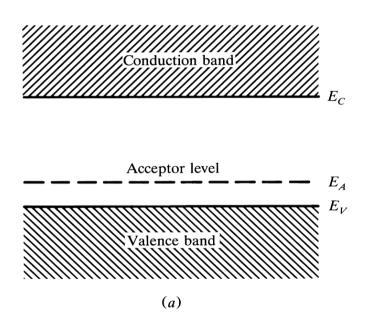


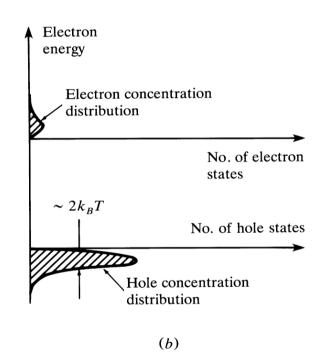






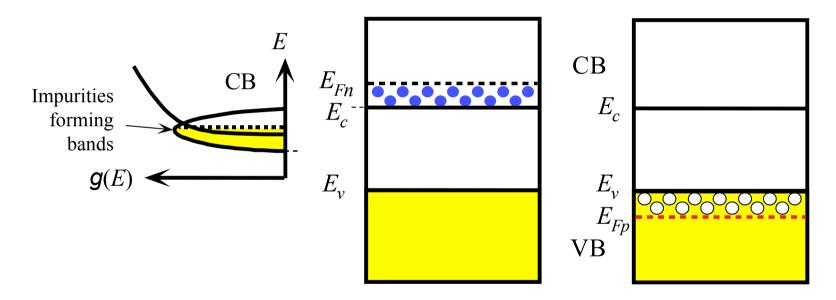
p-type material







☐ Heavily Doped Dopant States



Degenerated n-type semiconductor Large number of donors form a band that overlaps the CB

Degenerated p-type semiconductor



#### **Intrinsic semiconductor**

n = p and therefore:

$$2\left[\frac{2\pi m_e^* k_B T}{h^2}\right]^{3/2} e^{(E_F - E_c)/k_B T} = 2\left[\frac{2\pi m_h^* k_B T}{h^2}\right]^{3/2} e^{(E_v - E_F)/k_B T}$$

Here we can have a position of Fermi level in intrinsic SC:

$$E_F(T) = \frac{E_g}{2} + \frac{3}{4} k_B T \ln \left[ \frac{m_h^*}{m_e^*} \right]$$

So, it is almost in the middle of the bandgap



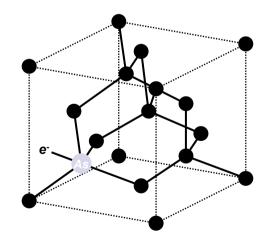
### **Extrinsic Semiconductors**

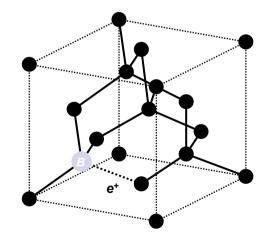
- \* Charge-neutrality equation
  - ⇒ Evaluation of carrier concentrations
- \* Fermi-level variation in extrinsic semiconductors
  - **⇒** Doping dependence
- \* Impurity statistics



## Charge-Neutrality Equation

- Previously we discussed how we may improve the electrical properties of semiconductors by the deliberate addition of impurities known as **DOPANTS** 
  - \* The doped semiconductor is referred to as **EXTRINSIC** since at room temperature its electrical characteristics are determined by the concentration of **DOPANTS** rather than the intrinsic carrier concentration  $n_i$ 
    - ⇒ Doping with DONORS increases the ELECTRON concentration and the resulting semiconductor is said to be *n*-TYPE while doping with ACCEPTORS increases the HOLE concentration and the resulting semiconductor is said to be *p*-TYPE





REPLACING A SILICON ATOM WITH ARSENIC YIELDS AN EXTRA ELECTRON

REPLACING A SILICON ATOM WITH BORON YIELDS AN EXTRA HOLE

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## Charge-Neutrality Equation

• The carrier concentrations in extrinsic semiconductors may be determined by introducing a CHARGE-NEUTRALITY equation which expresses the fact that the crystal as a whole must remain CHARGE NEUTRAL

$$p-n+N_{D}^{+}-N_{A}^{-}=0$$

- \* In this expression  $N_D^+$  and  $N_A^-$  are the concentrations of IONIZED donors and acceptors
  - ⇒ When an electron is liberated from a donor it is left with an equal POSITIVE charge while liberation of a hole from an acceptor leaves it NEGATIVELY charge
- \* Another important equation that will be used to determine the carrier concentrations in the extrinsic semiconductor is

$$np = 4\left[\frac{2\pi k_B T}{h^2}\right]^3 (m_e^* m_h^*)^{3/2} \exp\left[-\frac{E_g}{k_B T}\right]$$

⇒ Since the RHS of this equation is INDEPENDENT of the doping the PRODUCT of the electron and hole concentrations in a non-degenerate semiconductor is similarly INDEPENDENT of the doping concentration

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#### Fermi-Level Variation in Extrinsic Semiconductors

In doped semiconductors:

$$n = n_i \exp[(E_F - E_i)/k_B T]$$
$$p = n_i \exp[(E_i - E_F)/k_B T]$$

\* By REARRANGING these equations we obtain expressions relating the Fermi level to electron and hole concentrations in doped semiconductors:

$$E_F - E_i = k_B T \ln \left[ \frac{n}{n_i} \right] = -k_B T \ln \left[ \frac{p}{n_i} \right]$$

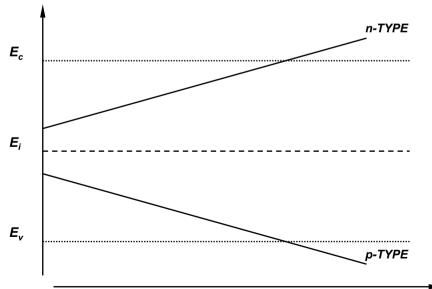


#### Fermi-Level Variation in Extrinsic Semiconductors

 For HEAVY doping and FULL ionization of dopants the position of the Fermi energy in the gap varies as

$$E_F - E_i = k_B T \ln \left[ \frac{N_D}{n_i} \right], \quad N_D > N_A >> n_i$$

$$E_F - E_i = -k_B T \ln \left[ \frac{N_A}{n_i} \right], \quad N_A > N_D >> n_i$$



- SCHEMATIC ILLUSTRATION INDICATING THE FERMI LEVEL VARIATION IN A SEMICONDUCTOR WITH n- AND p-TYPE DOPING
- THE DOPING RANGE ILLUSTRATED HERE CORRESPONDS TO THE REGIME OF HEAVY DOPING WHERE THE DOPANT DENSITIES SIGNIFICANTLY EXCEED n;
- NOTE THAT FOR SUFFICIENTLY HIGH DOPING DENSITIES THE FERMI LEVEL MOVES CLOSE TO EITHER BAND EDGE SO THAT THE SEMICONDUCTOR BECOMES DEGENERATE
- THE TERM "DEGENERATE" IS THEREFORE OFTEN USED TO INDICATE A VERY-HEAVILY DOPED SEMICONDUCTOR



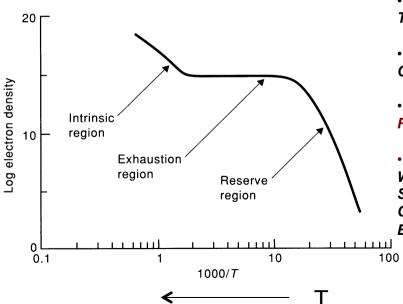
## Impurity Statistics

• At HIGH temperatures  $N_c >> N_D$  and

$$n = N_D$$

\* In this case the Fermi energy will lie well BELOW the donor energy as ALL the donors are ionized and the density is therefore relatively CONSTANT

\* The variation of the electron density with temperature in a typical *n*-type semiconductor is shown schematically in the figure below



- VARIATION OF THE ELECTRON DENSITY WITH TEMPERATURE IN A TYPICAL n-TYPE SEMICONDUCTOR
- THE REGION WHERE THE DENSITY IS INDEPENDENT OF TEMPERATURE IS KNOWN AS THE EXHAUSTION REGION
- AT LOWER TEMPERATURES THE DONORS BEGIN TO FREEZE-OUT CAUSING A REDUCTION OF THE ELECTRON CONCENTRATION
- THIS LOW-TEMPERATURE REGION IS OFTEN REFERRED TO AS THE RESERVE REGION
- AT HIGHER TEMPERATURES THAN WE HAVE CONSIDERED HERE WE MAY NO LONGER NEGLECT THE HOLE CONCENTRATION AND THE SEMICONDUCTOR REVERTS TO INTRINSIC-LIKE BEHAVIOR WHERE CARRIERS FROM THE BULK CRYSTAL SWAMP ELECTRONS PROVIDED BY THE DONORS